

Condensed Matter Physics

PHY 5111

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The Free Fermi Gas and Single Electron Model

Condensed Matter Physics – Michael P. Marder

Chapter 6



The Hamiltonian

Much of condensed matter physics lies within a Hamiltonian that one easily can write down in a single line. It is

$$\hat{\mathcal{H}} = \sum_l \frac{\hat{p}_l^2}{2M_l} + \frac{1}{2} \sum_{l \neq l'} \frac{q_l q_{l'}}{|\hat{R}_l - \hat{R}_{l'}|}.$$



The single-electron model

$$\sum_{l=1}^N \left(\frac{-\hbar^2 \nabla_l^2}{2m} + U(\vec{r}_l) \right) \Psi(\vec{r}_1 \dots \vec{r}_N) = \mathcal{E} \Psi(\vec{r}_1 \dots \vec{r}_N)$$

$$\left(\frac{-\hbar^2 \nabla^2}{2m} + U(\vec{r}) \right) \psi_l(\vec{r}) = \mathcal{E}_l \psi_l(\vec{r})$$



The free Fermi gas

$$\frac{-\hbar^2}{2m} \sum_{l=1}^N \nabla_l^2 \Psi(\vec{r}_1 \dots \vec{r}_N) = \mathcal{E} \Psi(\vec{r}_1 \dots \vec{r}_N)$$

To simplify further we impose periodic boundary conditions

$$\Psi(x_1 + L, y_1, z_1 \dots, z_N) = \Psi(x_1, y_1, z_1 \dots, z_N)$$

$$\Psi(x_1, y_1 + L, z_1 \dots, z_N) = \Psi(x_1, y_1, z_1 \dots, z_N)$$

⋮
⋮
⋮



One Free Fermion

$$\psi_{\vec{k}} = \frac{1}{\sqrt{\mathcal{V}}} e^{i\vec{k} \cdot \vec{r}}$$

$$L^3 = \mathcal{V}$$

$$\vec{k} = \frac{2\pi}{L} (l_x, l_y, l_z)$$

$l_x, l_y,$ and l_z are integers ranging from $-\infty$ to ∞

$$\epsilon_{\vec{k}}^0 = \frac{\hbar^2 k^2}{2m}$$



Many Free Fermions

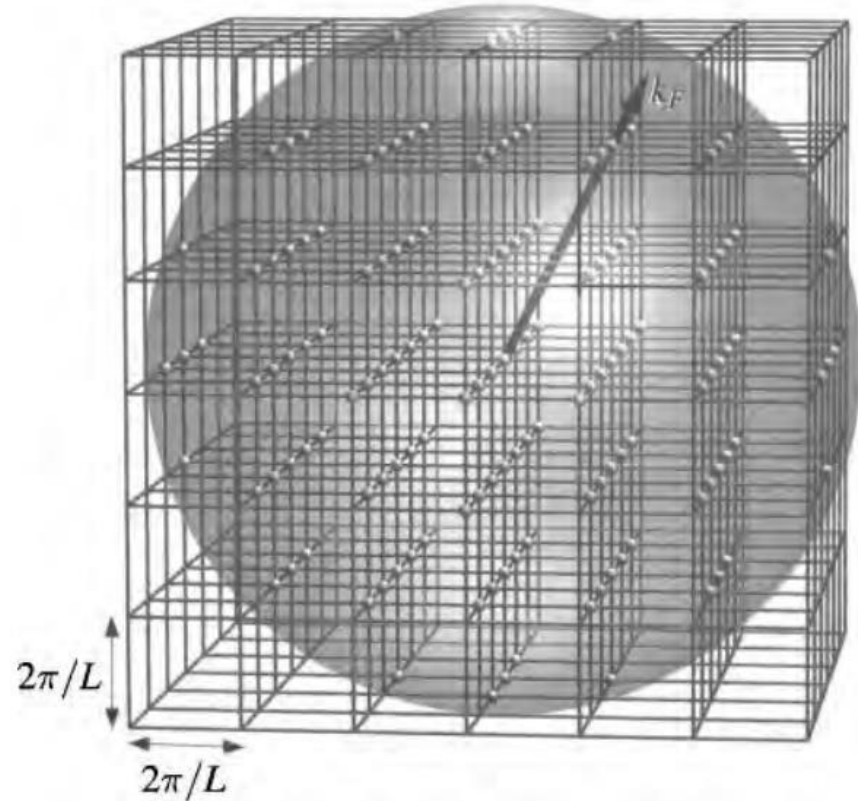
The ground state of electrons obeying free Fermi gas assumption is constructed from products of the one-electron wave functions. The Pauli exclusion principle forbids any given state from being occupied more than once, and therefore any given state indexed by \vec{k} is able to host no more than two electrons, one for each value of spin.



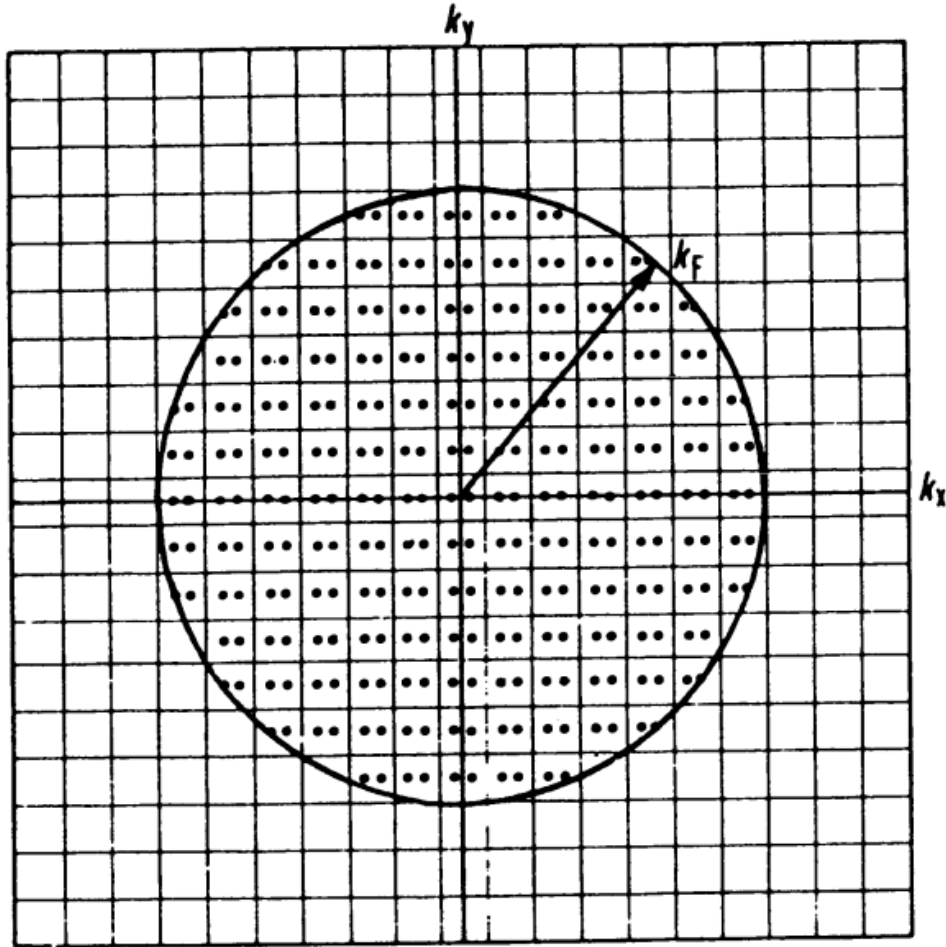
Many Free Fermions

$$\vec{k} = \frac{2\pi}{L} (l_x, l_y, l_z)$$

$l_x, l_y,$ and l_z are integers ranging from $-\infty$ to ∞



Density of States



Volume occupied by 2 states =

Total number of states =

States per unit volume =

Density of States

$$D_{\vec{k}} = 2 \frac{1}{(2\pi)^3}$$

For each wave vector Pauli's exclusion principle allows two electrons, one with spin up and the other with spin down.

$$\int [d\vec{k}] \equiv \frac{2}{V} \sum_{\vec{k}} = \int d\vec{k} D_{\vec{k}} = \frac{2}{(2\pi)^3} \int d\vec{k}$$



Energy Density of States

$$D(\mathcal{E}) = \int [d\vec{k}] \delta(\mathcal{E} - \mathcal{E}_{\vec{k}})$$

The units of densities of states are able to change without much warning. Often they are expressed in units of $1/[\text{eV atom}]$, which means they are related to the function defined by above equation by a factor of density n .



Energy Density of States

$$\begin{aligned} D(\mathcal{E}) &= \int [d\vec{k}] \delta(\mathcal{E} - \mathcal{E}_{\vec{k}}^0) \\ &= 4\pi \frac{2}{(2\pi)^3} \int_0^\infty dk k^2 \delta(\mathcal{E} - \mathcal{E}_{\vec{k}}^0) \\ &= \frac{1}{\pi^2} \int_0^\infty \frac{d\mathcal{E}^0}{|d\mathcal{E}^0/dk|} \frac{2m\mathcal{E}^0}{\hbar^2} \delta(\mathcal{E} - \mathcal{E}^0) \\ &= \frac{m}{\hbar^3 \pi^2} \sqrt{2m\mathcal{E}} \\ &= 6.812 \cdot 10^{21} \sqrt{\mathcal{E}/\text{eV}} \text{ eV}^{-1} \text{ cm}^{-3}. \end{aligned}$$

For the free
Fermi gas



Electron density

The number of electrons that can fit into a sphere of radius k_F is

$$N = \sum_{\vec{k}\sigma} f_{\vec{k}}$$
$$= \mathcal{V} \int [d\vec{k}] f_{\vec{k}},$$

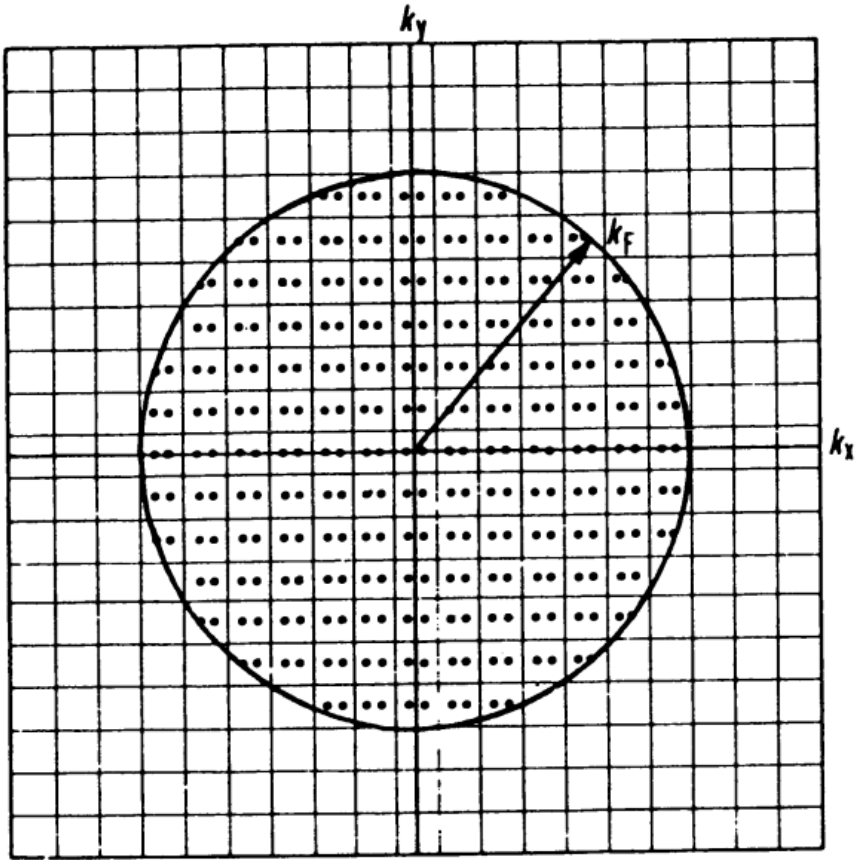
Because $f_{\vec{k}}$ is 1 only if the state is occupied.

$$= \frac{\mathcal{V}}{4\pi^3} \frac{4\pi}{3} k_F^3 = \frac{\mathcal{V} k_F^3}{3\pi^2},$$

$$n = N/\mathcal{V}$$



Electron density



$$4\pi \int_0^{k_F} g(k) k^2 dk = \frac{N}{V_g} = n$$

$$k_F = (3\pi^2 n)^{1/3}, \quad E_F = \frac{\hbar^2}{2m} (3\pi^2 n)^{2/3}$$

$$k_F = (3\pi^2 n)^{1/3} = 3.09 [n \cdot \text{\AA}^3]^{1/3} \text{\AA}^{-1}$$

$$\mathcal{E}_F = \frac{\hbar^2 k_F^2}{2m} = 36.46 [n \cdot \text{\AA}^3]^{2/3} \text{eV}$$

For the free
Fermi gas

$$v_F = \hbar k_F / m = 3.58 [n \cdot \text{\AA}^3]^{1/3} \cdot 10^8 \text{ cm s}^{-1}$$

$$D(\mathcal{E}_F) = \frac{3}{2} \frac{n}{\mathcal{E}_F} = 4.11 \cdot 10^{-2} [n \cdot \text{\AA}^3] \text{eV}^{-1} \text{\AA}^{-3}$$



Grand Partition Function

$$Z_{\text{gr}} = \sum_{\text{states}} e^{\beta(\mu N - \mathcal{E})}$$

$$= \sum_{n_1=0}^1 \sum_{n_2=0}^1 \sum_{n_3=0}^1 \dots e^{\beta \sum_l n_l (\mu - \mathcal{E}_l)}$$

$$= \prod_l \left\{ \sum_{n_l=0}^1 e^{\beta n_l [\mu - \mathcal{E}_l]} \right\}$$

$$= \prod_l \left[1 + e^{\beta [\mu - \mathcal{E}_l]} \right].$$

$$\sum_{n_1=0}^N \sum_{n_2=0}^N \dots \sum_{n_M=0}^N \prod_{l=1}^M A_{n_l} = \prod_{l=1}^M \left\{ \sum_{n_l=0}^N A_{n_l} \right\}$$



Grand Potential

$$\begin{aligned}\Pi &\equiv -k_B T \ln Z_{\text{gr}} \\ &= -k_B T \sum_l \ln \left[1 + e^{\beta[\mu - \varepsilon_l]} \right] . \\ &= -k_B T \mathcal{V} \int d\varepsilon D(\varepsilon) \ln \left[1 + e^{\beta[\mu - \varepsilon]} \right] .\end{aligned}$$



Fermi Distribution Function

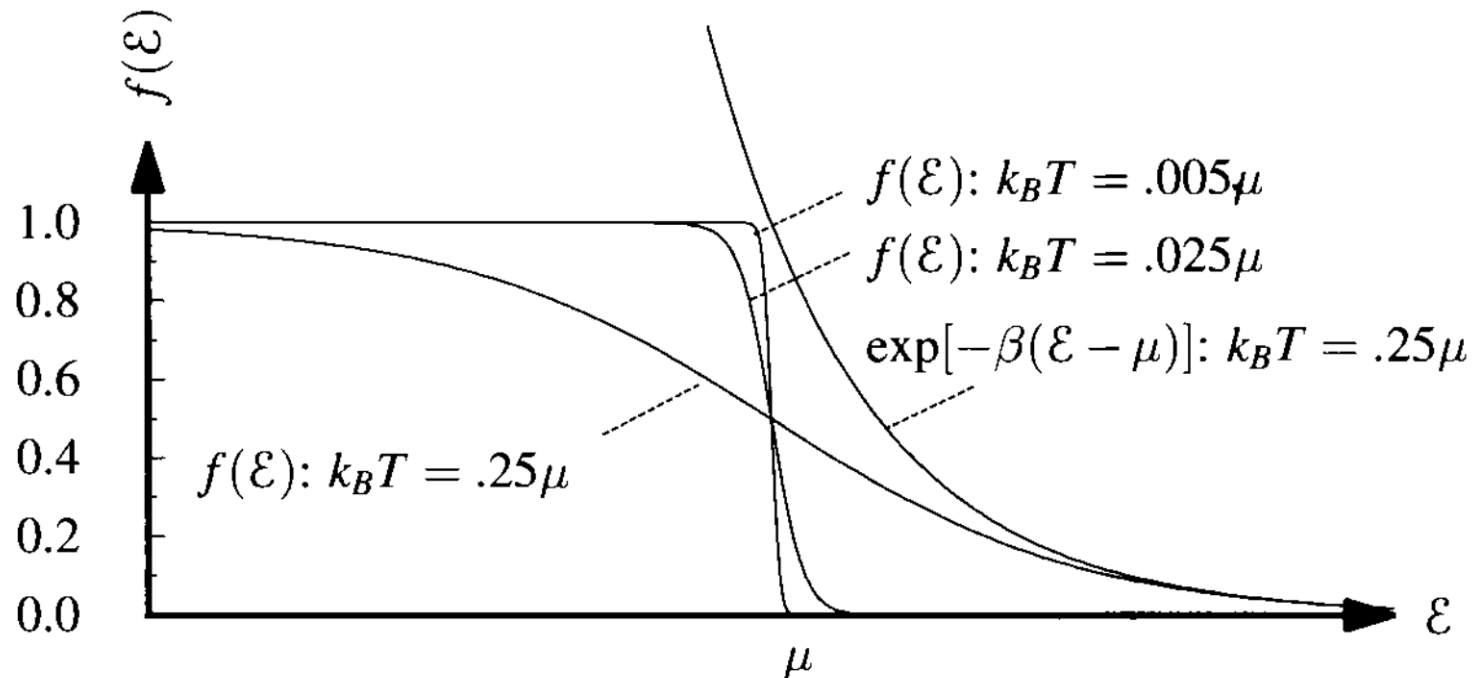
$$\begin{aligned} N &= -\frac{\partial \Pi}{\partial \mu} \\ &= \mathcal{V} \int d\varepsilon' D(\varepsilon') \frac{e^{\beta\mu - \beta\varepsilon'}}{1 + e^{\beta\mu - \beta\varepsilon'}} \\ \Rightarrow n &= \frac{N}{\mathcal{V}} = \int d\varepsilon' D(\varepsilon') f(\varepsilon'), \end{aligned}$$

$$f(\varepsilon) = \frac{1}{e^{\beta(\varepsilon - \mu)} + 1}$$

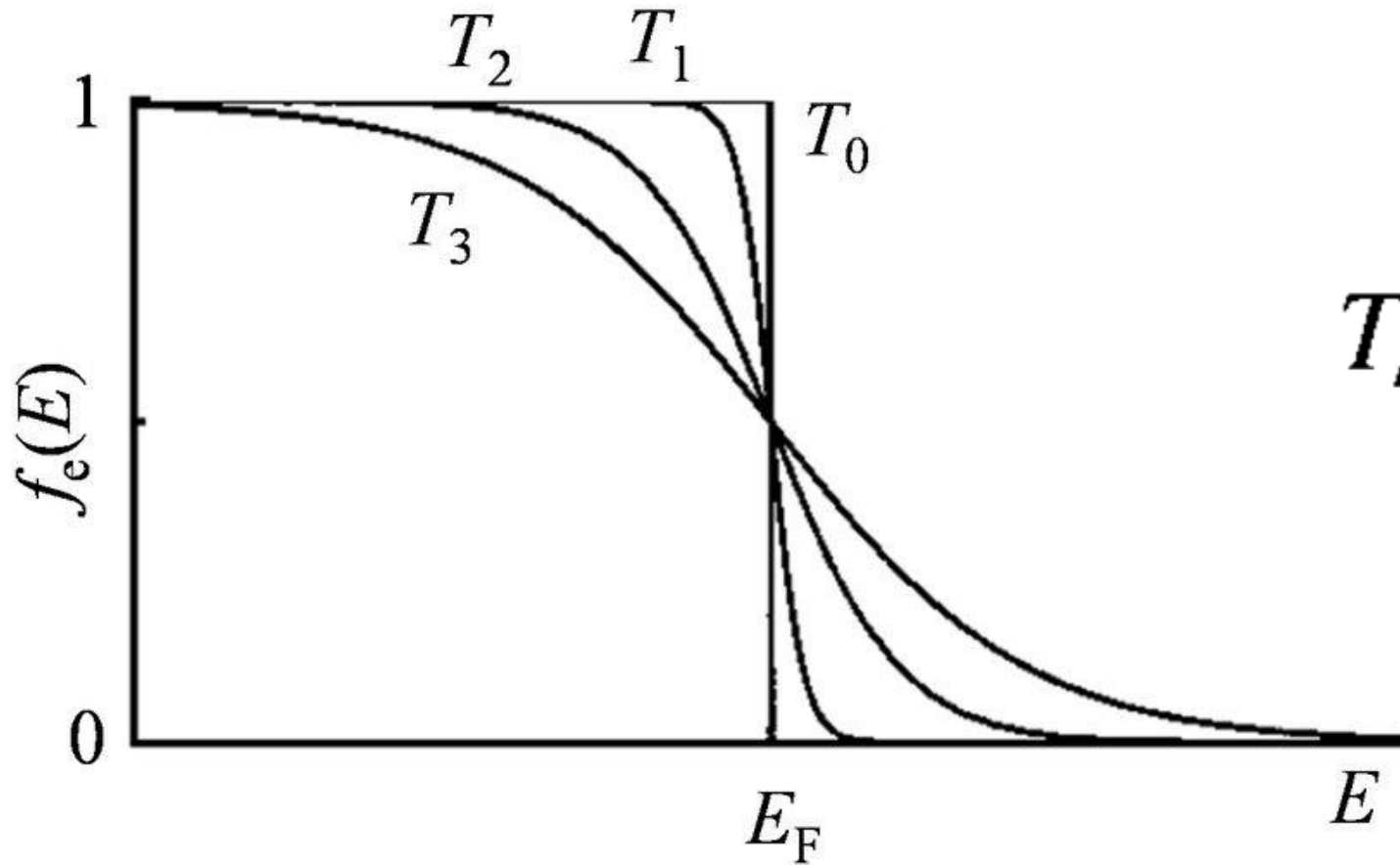


Fermi Distribution Function

$$f(\mathcal{E}) = \frac{1}{e^{\beta(\mathcal{E}-\mu)} + 1}$$

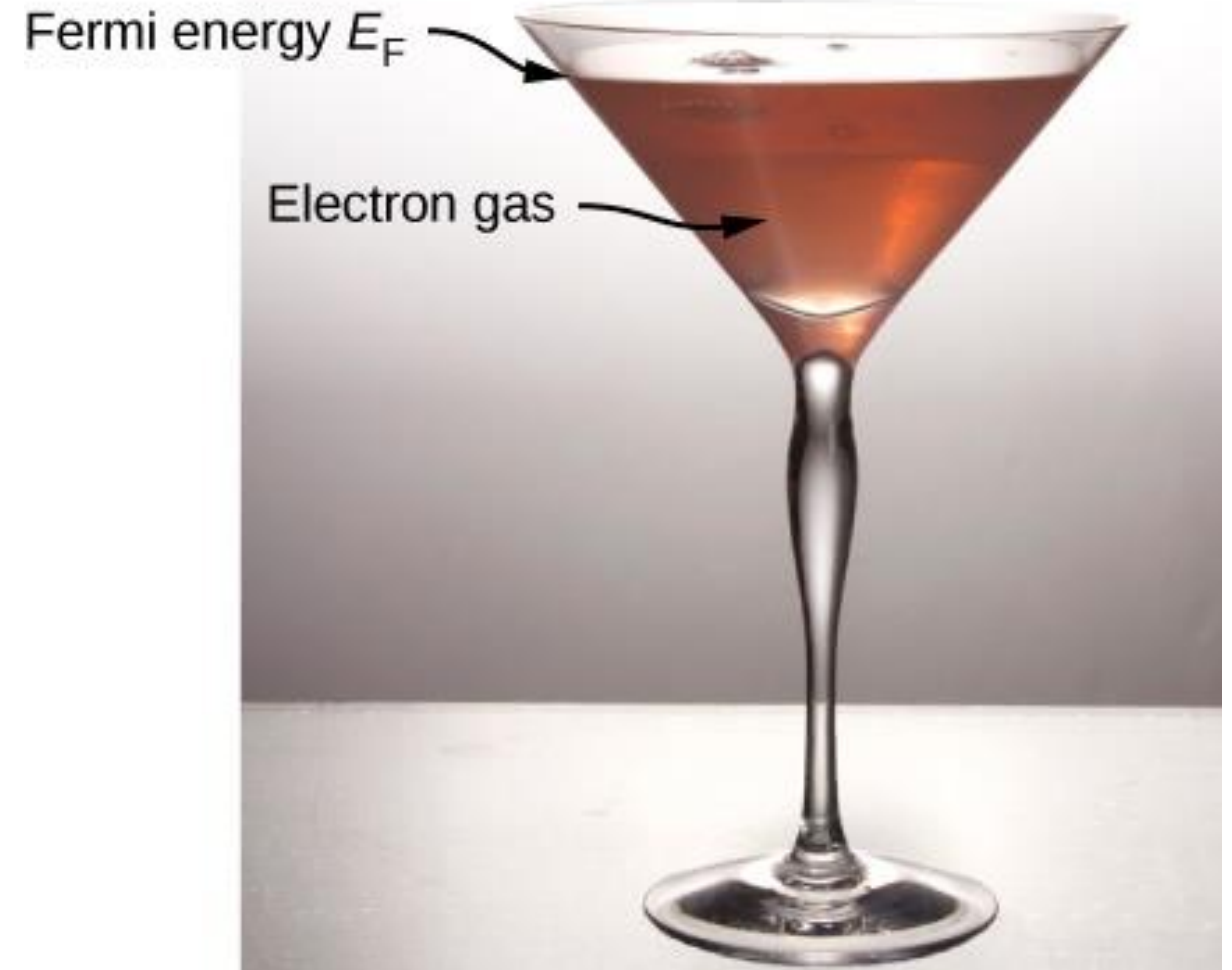
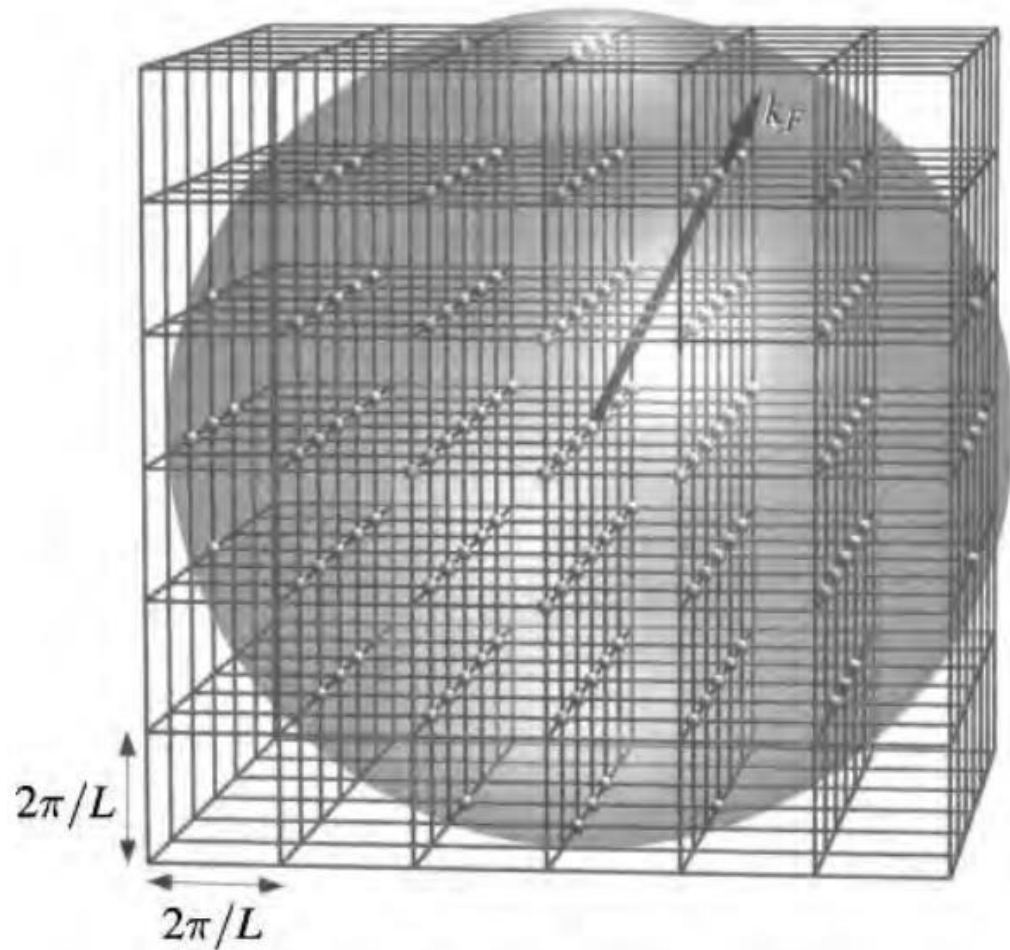


Fermi Distribution Function



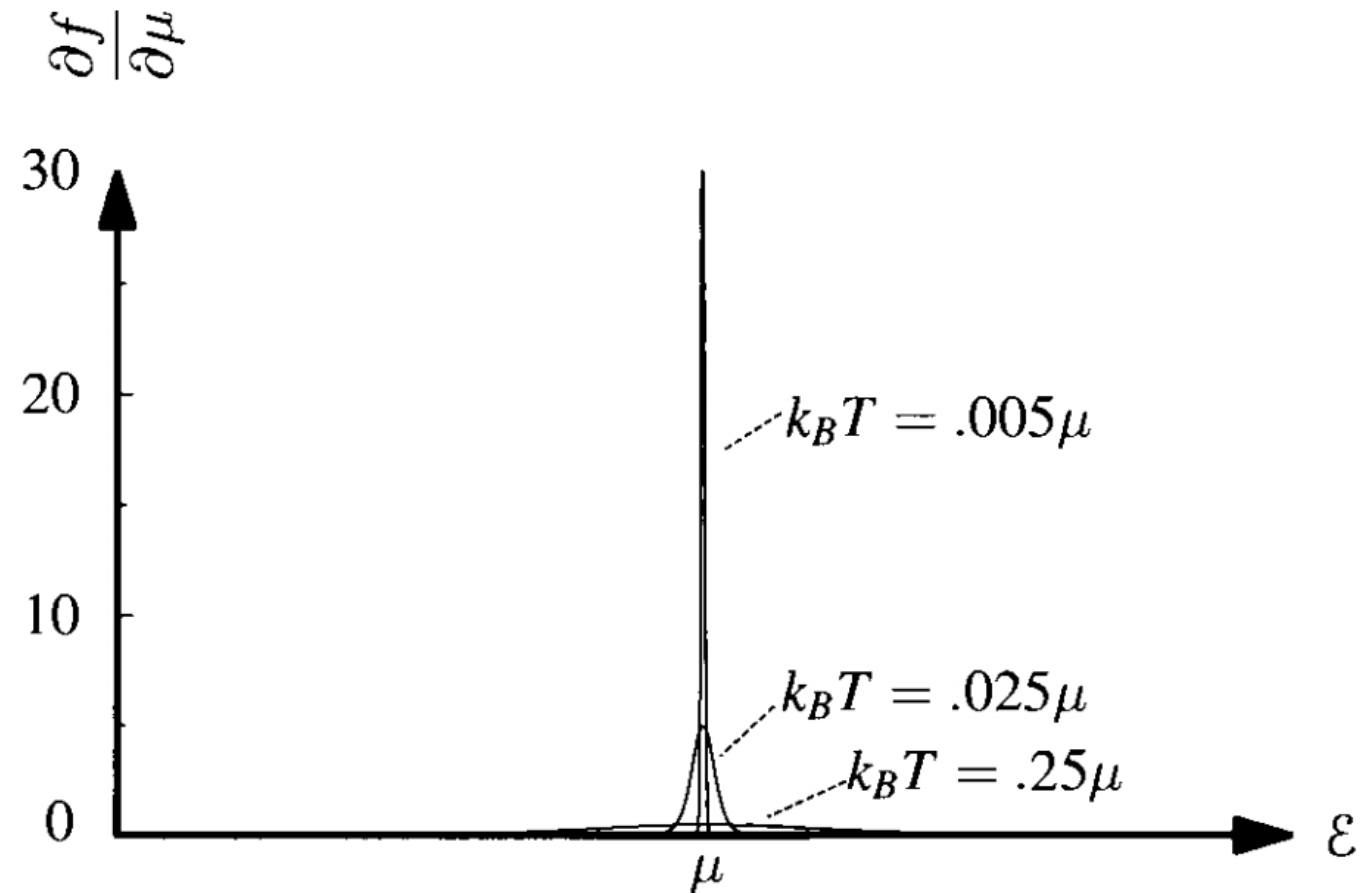
$$T_F = \mathcal{E}_F / k_B,$$

Fermi Surface



Element	Z	n (10^{22} cm^{-3})	k_F (10^8 cm^{-1})	\mathcal{E}_F (eV)	T_F (10^4 K)	v_F (10^8 cm s^{-1})	r_s/a_0
Li	1	4.60	1.11	4.68	5.43	1.28	3.27
Na	1	2.54	0.91	3.15	3.66	1.05	3.99
K	1	1.32	0.73	2.04	2.37	0.85	4.95
Rb	1	1.08	0.68	1.78	2.06	0.79	5.30
Cs	1	0.85	0.63	1.52	1.76	0.73	5.75
Cu	1	8.49	1.36	7.04	8.17	1.57	2.67
Ag	1	5.86	1.20	5.50	6.38	1.39	3.02
Au	1	5.90	1.20	5.53	6.42	1.39	3.01
Be	2	24.72	1.94	14.36	16.67	2.25	1.87
Mg	2	8.62	1.37	7.11	8.26	1.58	2.65
Ca	2	4.66	1.11	4.72	5.48	1.29	3.26
Sr	2	3.49	1.01	3.89	4.52	1.17	3.59

The derivative of the Fermi function



Specific Heat of Noninteracting Electrons at Low T

$$c_V = \frac{1}{V} \frac{\partial \mathcal{E}}{\partial T} \Big|_{N, V}$$

$$c_V = \frac{\pi^2}{3} k_B^2 T D(\mathcal{E}_F)$$

