#### Dr Mohammad Abdur Rashid



Jashore University of Science and Technology **Dr Rashid, 2020**

#### Schrödinger Equation

$$
i\hbar \frac{\partial}{\partial t} \Psi(x,t) = \hat{H} \Psi(x,t)
$$

For a particle quantum mechanical of mass  $m$  moving in a potential  $V(x,t)$  the Hamiltonian operator is

$$
\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x, t)
$$



#### Hermitian operator

#### An operator  $\hat{\Omega}$ , which corresponds to a physical observable  $\Omega$ , is said to be Hermitian if

$$
\int \Phi^* \hat{\Omega} \Psi \, dx = \int (\hat{\Omega} \Phi)^* \Psi \, dx
$$

**Examples:** position operator  $\hat{x}$ , momentum operator  $\hat{p}$ , Hamiltonian operator  $\hat{H}$ , etc.



#### Expectation value

The expectation value  $\langle \Omega \rangle$  of any operator  $\Omega$ :

$$
\langle \Omega \rangle = \int_{-\infty}^{\infty} \Psi^*(x, t) \, \hat{\Omega} \Psi(x, t) \, dx
$$

The physical meaning of this would be if we consider many copies of the identical system, and measure  $\Omega$  at a time t in all of them, then the average value recorded will converge to  $\langle \Omega \rangle$  as the number of systems and measurements approaches infinity.



$$
i\hbar \frac{d}{dt}\langle \Omega \rangle = i\hbar \frac{d}{dt} \int_{-\infty}^{\infty} \Psi^*(x, t) \hat{\Omega} \Psi(x, t) dx
$$

$$
= i\hbar \int_{-\infty}^{\infty} \left( \frac{\partial \Psi^*}{\partial t} \hat{\Omega} \Psi + \Psi^* \hat{\Omega} \frac{\partial \Psi}{\partial t} \right) dx
$$

$$
\boxed{ i\hbar \frac{\partial \Psi}{\partial t} = \hat{H} \Psi}
$$

$$
\boxed{-i\hbar\frac{\partial\Psi^*}{\partial t}=\big(\hat{H}\Psi\big)^*}
$$



$$
i\hbar \frac{d}{dt}\langle \Omega \rangle = \int_{-\infty}^{\infty} \left\{ -(\hat{H}\Psi)^* \hat{\Omega}\Psi + \Psi^* \hat{\Omega}(\hat{H}\Psi) \right\} dx
$$
  
= 
$$
\int_{-\infty}^{\infty} \left\{ \Psi^* \hat{\Omega}(\hat{H}\Psi) - (\hat{H}\Psi)^* \hat{\Omega}\Psi \right\} dx
$$

Hermitian operator:

$$
\int \Phi^* \hat{\Omega} \Psi \, dx = \int (\hat{\Omega} \Phi)^* \Psi \, dx
$$



$$
i\hbar \frac{d}{dt}\langle \Omega \rangle = \int_{-\infty}^{\infty} (\Psi^* \hat{\Omega} \hat{H} \Psi - \Psi^* \hat{H} \hat{\Omega} \Psi) dx
$$
  

$$
= \int_{-\infty}^{\infty} \Psi^* (\hat{\Omega} \hat{H} - \hat{H} \hat{\Omega}) \Psi dx
$$
  

$$
= \int_{-\infty}^{\infty} \Psi^* [\hat{\Omega}, \hat{H}] \Psi dx
$$

Commutator of two operators:

$$
[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A}
$$



We have proven that for operators  $\hat{\Omega}$  that do not explicitly depend on time,

$$
i\hbar\frac{\mathrm{d}}{\mathrm{d}t}\langle\Omega\rangle=\left\langle[\hat{\Omega},\,\hat{H}]\right\rangle
$$

Therefore, for an operation which commute with the Hamiltonian operator  $\hat{H}$  the expectation value will not change over time.



#### **Example:** Consider the momentum operator  $\hat{p}$  for a free particle.

# For a free particle the Hamiltonian operator:  $\hat{H} = \frac{\hat{p}^2}{2m}$



Hence the momentum operator  $\hat{p}$  commute with H:

$$
\begin{array}{lcl} [\hat{p},\,\hat{H}] & = & [\hat{p},\,\frac{\hat{p}^2}{2m}]=\frac{1}{2m}[\hat{p},\,\hat{p}^2] \\ & = & \frac{1}{2m}\Big([\hat{p},\,\hat{p}]\hat{p}+\hat{p}[\hat{p},\,\hat{p}]\Big) = 0 \end{array}
$$

$$
\left[ [\hat{A}, \hat{B}\hat{C}] = [\hat{A}, \hat{B}]\hat{C} + \hat{B}[\hat{A}, \hat{C}] \right]
$$



 $i\hslash \frac{\mathrm{d}}{\mathrm{d}t}\langle p\rangle=\left\langle \left[\hat{p},\,\hat{H}\right]\right\rangle =0$ 

For a free particle the expectation value of momentum does not change over time.



### Thank You

#### **To receive notification of new video please subscribe to our channel.**

You may also let us know your comments.



Jashore University of Science and Technology **Dr Rashid, 2020**