

Quantum Mechanics

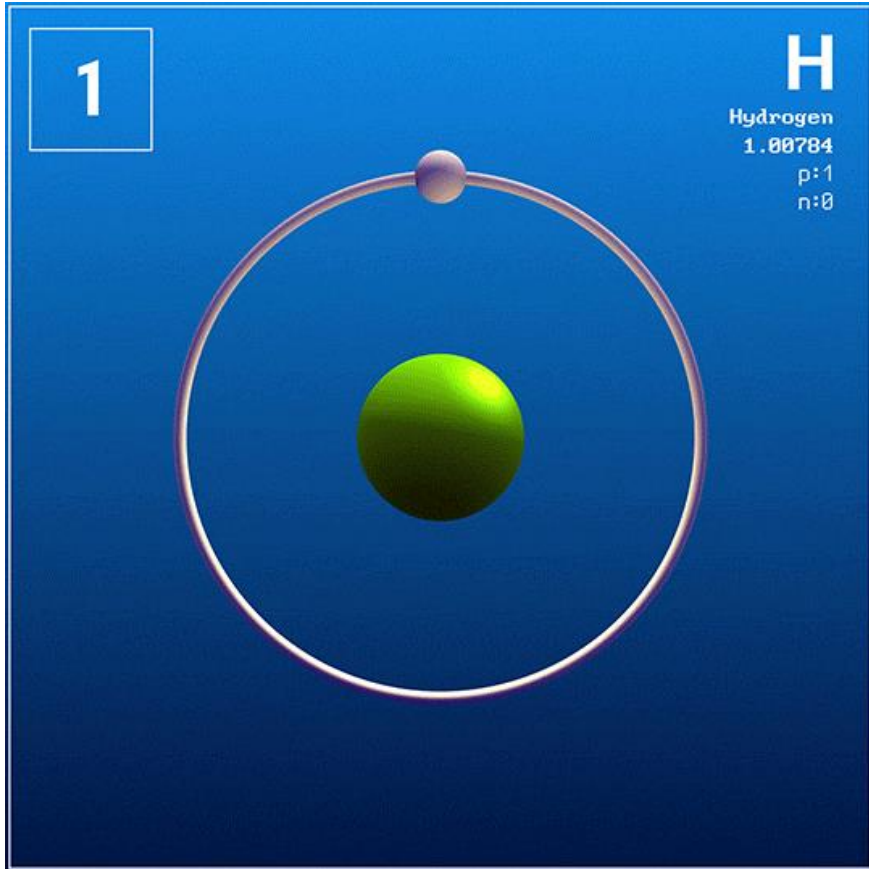
The Hydrogen Like Atom

Hydrogenic wave function

Dr Mohammad Abdur Rashid



Hydrogen atom



Wikipedia

| Particle | Position | Momentum |
|----------|----------------|----------------|
| Electron | \mathbf{r}_1 | \mathbf{p}_1 |
| Proton | \mathbf{r}_2 | \mathbf{p}_2 |

$$[(\hat{\mathbf{r}}_1)_i, (\hat{\mathbf{p}}_1)_j] = i\hbar\delta_{ij}$$

$$[(\hat{\mathbf{r}}_2)_i, (\hat{\mathbf{p}}_2)_j] = i\hbar\delta_{ij}$$

The Hydrogenic Atom

A hydrogen atom or a hydrogen like atom (He^+ , Li^{2+} , Be^{+3} , etc.) consists of an atomic nucleus of charge Ze and an electron of charge $-e$. Their mutual interaction is given by the Coulomb potential

$$V(|\mathbf{r}_1 - \mathbf{r}_2|) = -\frac{1}{4\pi\epsilon_0} \frac{Ze^2}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

where $\mathbf{r}_1 = \mathbf{r}_1(x_1, y_1, z_1)$ and $\mathbf{r}_2 = \mathbf{r}_2(x_2, y_2, z_2)$ are the electron and nucleus position vectors, respectively.



The Schrödinger equation

The time-independent Schrödinger equation for the system is given by

$$\left\{ -\frac{\hbar^2}{2m_1} \nabla_1^2 - \frac{\hbar^2}{2m_2} \nabla_2^2 + V(|\mathbf{r}_1 - \mathbf{r}_2|) \right\} \Psi(\mathbf{r}_1, \mathbf{r}_2) = E_{\text{tot}} \Psi(\mathbf{r}_1, \mathbf{r}_2),$$

where m_1 and m_2 are the masses of electron and nucleus.

$$\nabla_i^2 = \frac{\partial^2}{\partial x_i^2} + \frac{\partial^2}{\partial y_i^2} + \frac{\partial^2}{\partial z_i^2}; \quad i = 1, 2$$



Change of variables

$$\Psi(\mathbf{r}_1, \mathbf{r}_2) = \Psi(\mathbf{R}, \mathbf{r})$$

$$\mathbf{r} = \mathbf{r}_1 - \mathbf{r}_2$$

$$\mathbf{R} = \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2}$$

$$\mathbf{p} = \mu \left(\frac{\mathbf{p}_1}{m_1} - \frac{\mathbf{p}_2}{m_2} \right)$$

$$\mathbf{P} = \mathbf{p}_1 + \mathbf{p}_2.$$

$$[(\hat{\mathbf{r}}_1)_i, (\hat{\mathbf{p}}_1)_j] = i\hbar\delta_{ij}$$

$$[(\hat{\mathbf{r}}_2)_i, (\hat{\mathbf{p}}_2)_j] = i\hbar\delta_{ij}$$

$$[\hat{\mathbf{r}}_i, \hat{\mathbf{p}}_j] = i\hbar\delta_{ij},$$

$$[\hat{\mathbf{R}}_i, \hat{\mathbf{P}}_j] = i\hbar\delta_{ij}.$$



The Schrödinger equation in new variables

Since \mathbf{R} and \mathbf{r} are independent to each other the wave function $\Psi(\mathbf{R}, \mathbf{r})$ can be separated into a product of functions of the centre of mass coordinate \mathbf{R} and of relative coordinate \mathbf{r} as $\Psi(\mathbf{R}, \mathbf{r}) = \Phi(\mathbf{R})\psi(\mathbf{r})$. With this the Schrödinger equation can be written as

$$\left\{ -\frac{\hbar^2}{2M} \nabla_{\mathbf{R}}^2 - \frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right\} \Phi(\mathbf{R})\psi(\mathbf{r}) = E_{\text{tot}} \Phi(\mathbf{R})\psi(\mathbf{r})$$



The Schrödinger equation in new variables

$$\left\{ -\frac{\hbar^2}{2M} \nabla_R^2 - \frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right\} \Phi(\mathbf{R})\psi(\mathbf{r}) = E_{\text{tot}} \Phi(\mathbf{R})\psi(\mathbf{r})$$

$$-\frac{\hbar^2}{2M} \psi(\mathbf{r}) \nabla_R^2 \Phi(\mathbf{R}) + \Phi(\mathbf{R}) \left\{ -\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right\} \psi(\mathbf{r}) = E_{\text{tot}} \Phi(\mathbf{R})\psi(\mathbf{r})$$

$$-\frac{\hbar^2}{2M} \frac{1}{\Phi(\mathbf{R})} \nabla_R^2 \Phi(\mathbf{R}) + \frac{1}{\psi(\mathbf{r})} \left\{ -\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right\} \psi(\mathbf{r}) = E_{\text{tot}}.$$



Two separate Schrödinger equations

Thus, we have the following two separate equations

$$-\frac{\hbar^2}{2M} \nabla_{\mathbf{R}}^2 \Phi(\mathbf{R}) = E_{\text{CM}} \Phi(\mathbf{R})$$

$$\left\{ -\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \right\} \psi(\mathbf{r}) = E \psi(\mathbf{r})$$

with the condition $E_{\text{tot}} = E_{\text{CM}} + E$.

$$V(r) = -\frac{1}{4\pi\epsilon_0} \frac{Ze^2}{r}.$$



The center of mass equation

$$-\frac{\hbar^2}{2M} \nabla_{\mathbf{R}}^2 \Phi(\mathbf{R}) = E_{\text{CM}} \Phi(\mathbf{R})$$

The solution to this kind of equation has the form

$$\Phi(\mathbf{R}) = (2\pi)^{-3/2} e^{i\mathbf{k}\cdot\mathbf{R}}$$

where \mathbf{k} is the wave vector associated with the center of mass. The constant $E_{\text{CM}} = \hbar^2 k^2 / (2M)$ gives the kinetic energy of the center of mass in the laboratory system (the total mass M is located at the origin of the center of mass coordinate system).



The Hamiltonian in spherical polar coordinates

$$\begin{aligned} H &= -\frac{\hbar^2}{2\mu} \nabla^2 + V(r) \\ &= -\frac{\hbar^2}{2\mu} \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] + V(r) \\ &= -\frac{\hbar^2}{2\mu} \left[\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) - \frac{\mathbf{L}^2}{\hbar^2 r^2} \right] + V(r), \end{aligned}$$

where \mathbf{L}^2 is the square of the magnitude of the orbital angular momentum and defined as

$$\mathbf{L}^2 = -\hbar^2 \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right].$$



The time-independent Schrödinger equation

$$\left\{ -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{\mathbf{L}^2}{2\mu r^2} + V(r) \right\} \psi(\mathbf{r}) = E\psi(\mathbf{r}).$$

In order to simplify the solution of this equation we notice that \mathbf{L}^2 do not operate on the radial variable r . Since the spherical harmonics $Y_{lm}(\theta, \phi)$ are eigenfunctions of \mathbf{L}^2 we can look for solution of the Schrödinger equation having the separable form

$$\psi(\mathbf{r}) = \psi(r, \theta, \phi) = R_l(r)Y_{lm}(\theta, \phi)$$

where $R_l(r)$ is the radial function which remains to be found.



Spherical harmonics $Y_{lm}(\theta, \phi)$

$$\mathbf{L}^2 Y_{lm}(\theta, \phi) = \hbar^2 l(l+1) Y_{lm}(\theta, \phi)$$

$$Y_{lm}(\theta, \phi) = (-1)^m \sqrt{\left(\frac{2l+1}{4\pi}\right) \frac{(l-m)!}{(l+m)!}} P_l^m(\cos \theta) e^{im\phi}, \quad (m \geq 0).$$

Here $P_l^m(\cos \theta)$ is the associated Legendre functions. $m < 0$, we use

$$Y_{l,m}(\theta, \phi) = (-1)^m [Y_{l,-m}(\theta, \phi)]^*.$$

$$\int_0^{2\pi} d\phi \int_0^\pi d\theta \sin \theta Y_{l'm'}^*(\theta, \phi) Y_{lm}(\theta, \phi) = (Y_{l'm'}^*, Y_{lm}) = \delta_{l'l} \delta_{m'm}.$$



Solution of the Radial Equation

$$\mathbf{L}^2 Y_{lm}(\theta, \phi) = \hbar^2 l(l+1) Y_{lm}(\theta, \phi)$$

$$\left\{ -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{\hbar^2 l(l+1)}{2\mu r^2} + V(r) \right\} R_l(r) = E R_l(r)$$

$$\frac{d^2 R_l(r)}{dr^2} + \frac{2}{r} \frac{dR_l(r)}{dr} + \left[\frac{2\mu}{\hbar^2} E - \frac{l(l+1)}{r^2} + \frac{2\mu}{\hbar^2} \left(\frac{Ze^2}{4\pi\epsilon_0} \right) \frac{1}{r} \right] R_l(r) = 0$$



Asymptotic solution of the Radial Equation

Asymptotic solution: $r \rightarrow \infty$

$$\frac{d^2 R_l(r)}{dr^2} \approx -\frac{2\mu E}{\hbar^2} R_l(r) = \frac{2\mu|E|}{\hbar^2} R_l(r)$$

having noted that the energy E is negative for bound states.

$$R_l(r) = Ae^{-\sqrt{2\mu|E|/\hbar^2} r} + Be^{\sqrt{2\mu|E|/\hbar^2} r}$$

where A and B are constants to be determined.



Asymptotic solution of the Radial Equation

$$R_l(r) = Ae^{-\sqrt{2\mu|E|/\hbar^2} r} + Be^{\sqrt{2\mu|E|/\hbar^2} r}$$

Choose the negative exponential ($B = 0$) and set

$$E = -\frac{\mu Z^2 e^4}{8\epsilon_0^2 h^2} = -\frac{\mu Z^2 e^4}{2(4\pi\epsilon_0)^2 \hbar^2},$$

the ground state energy in the Bohr theory (in center of mass system), we obtain

$$R_l(r) = Ae^{-Zr/a_\mu}$$



Asymptotic solution of the Radial Equation

$$R_l(r) = Ae^{-Zr/a_\mu}$$

where a_μ is the modified Bohr radius

$$a_\mu = \frac{(4\pi\epsilon_0)\hbar^2}{\mu e^2} = \frac{\epsilon h^2}{\pi\mu e^2} = \frac{m_1}{\mu} \frac{\epsilon h^2}{\pi m_1 e^2} = \frac{m_1}{\mu} a_0$$

with a_0 being the Bohr radius.

$$\int_0^\infty [R_{10}(r)]^2 r^2 dr = 1$$

$$\int_0^\infty r^n e^{-\alpha r} dr = n! \alpha^{-(n+1)}$$



Normalized radial function

$$R_{10}(r) = 2 \left(\frac{Z}{a_{\mu}} \right)^{3/2} e^{-Zr/a_{\mu}}$$

$$\psi(\mathbf{r}) = \psi(r, \theta, \phi) = R_l(r) Y_{lm}(\theta, \phi)$$

$$\mathbf{L}^2 Y_{lm}(\theta, \phi) = \hbar^2 l(l+1) Y_{lm}(\theta, \phi)$$



The wave function ψ_{1s}

$$R_{10}(r) = 2 \left(\frac{Z}{a_{\mu}} \right)^{3/2} e^{-Zr/a_{\mu}}$$

$$Y_{00}(\theta, \phi) = 1/\sqrt{4\pi}$$

$$\psi_{100}(r, \theta, \phi) = \psi_{100}(r) = \psi_{1s}(r) = \frac{1}{\sqrt{\pi}} \left(\frac{Z}{a_{\mu}} \right)^{3/2} e^{-Zr/a_{\mu}}$$

The wave function of the hydrogen atom in ground state is found by setting $Z = 1$ as

$$\psi_{1s}(r) = \left(\frac{1}{\pi^{1/3} a_{\mu}} \right)^{3/2} e^{-r/a_{\mu}}$$



General solution of the radial wave function

The normalized radial function for the bound state of hydrogenic atom has a rather complicated form which we give without proof:

$$R_{nl}(r) = - \left\{ \left(\frac{2Z}{na_{\mu}} \right)^3 \frac{(n-l-1)!}{2n[(n+1)!]^3} \right\}^{1/2} e^{-\rho/2} \rho^l L_{n+l}^{2l+1}(\rho)$$

with

$$\rho = \frac{2Z}{na_{\mu}} r, \quad a_{\mu} = \frac{(4\pi\epsilon_0)\hbar^2}{\mu e^2}.$$

Here L_{β}^{α} is an associated Laguerre polynomial.



Radial eigenfunctions of hydrogenic atom

$$R_{10}(r) = 2 \left(\frac{Z}{a_{\mu}} \right)^{3/2} e^{-Zr/a_{\mu}}$$

$$R_{20}(r) = 2 \left(\frac{Z}{2a_{\mu}} \right)^{3/2} \left(1 - \frac{Zr}{2a_{\mu}} \right) e^{-Zr/2a_{\mu}}$$

$$R_{21}(r) = \frac{1}{\sqrt{3}} \left(\frac{Z}{2a_{\mu}} \right)^{3/2} \left(\frac{Zr}{a_{\mu}} \right) e^{-Zr/2a_{\mu}}$$



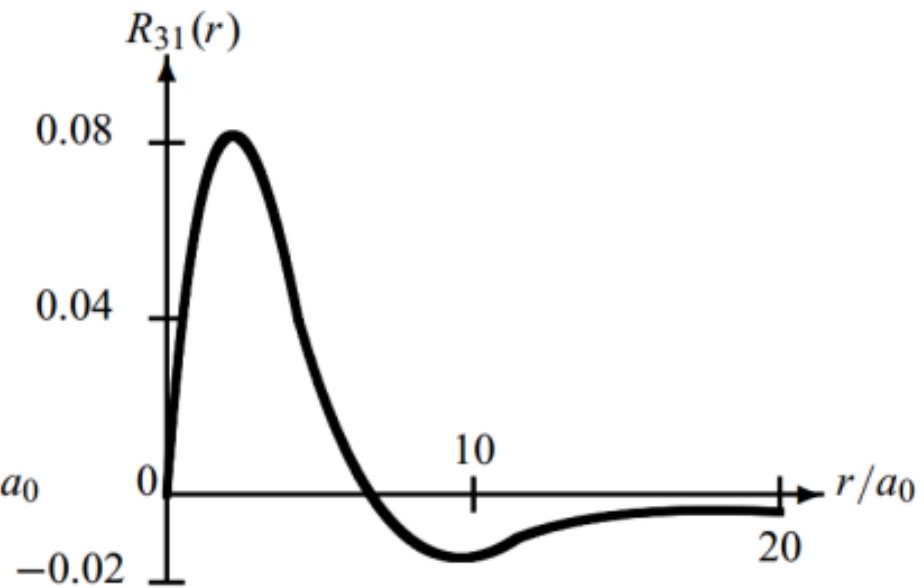
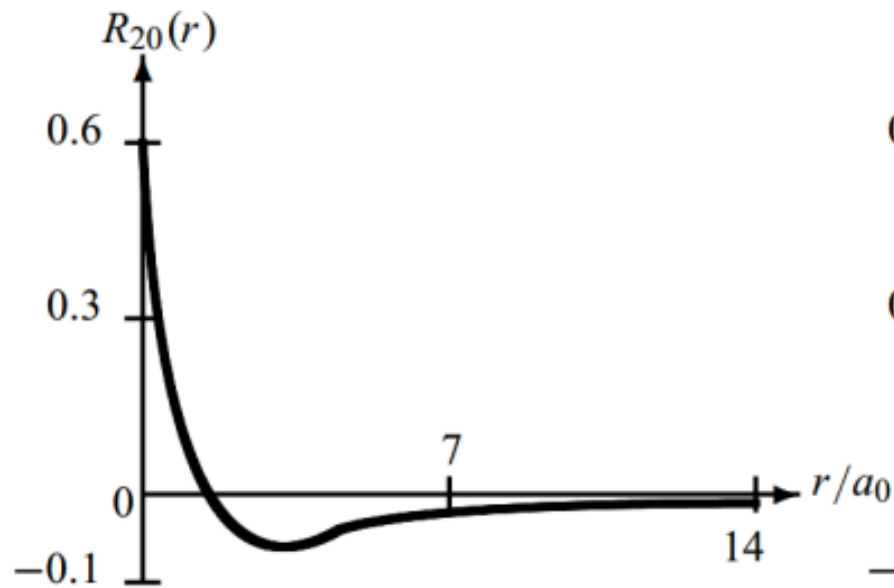
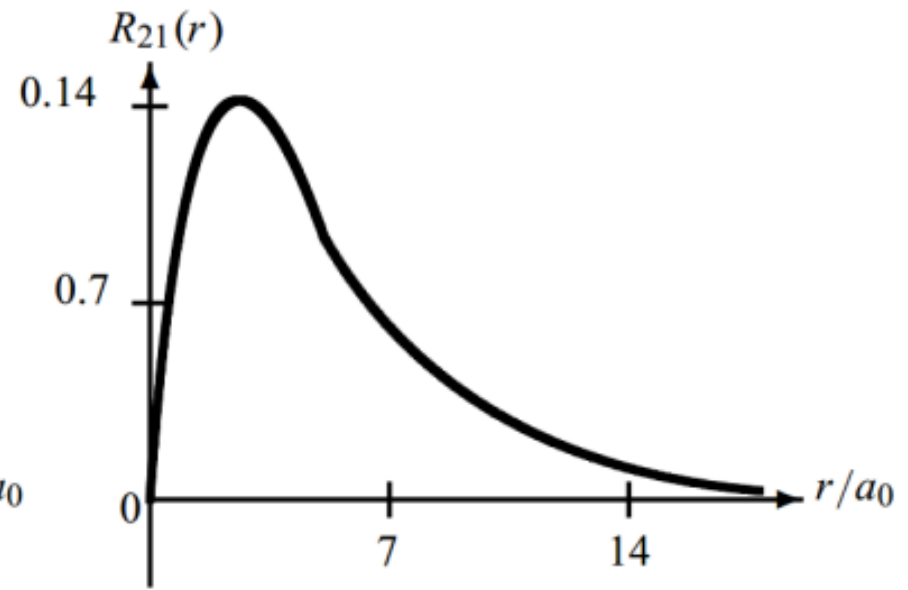
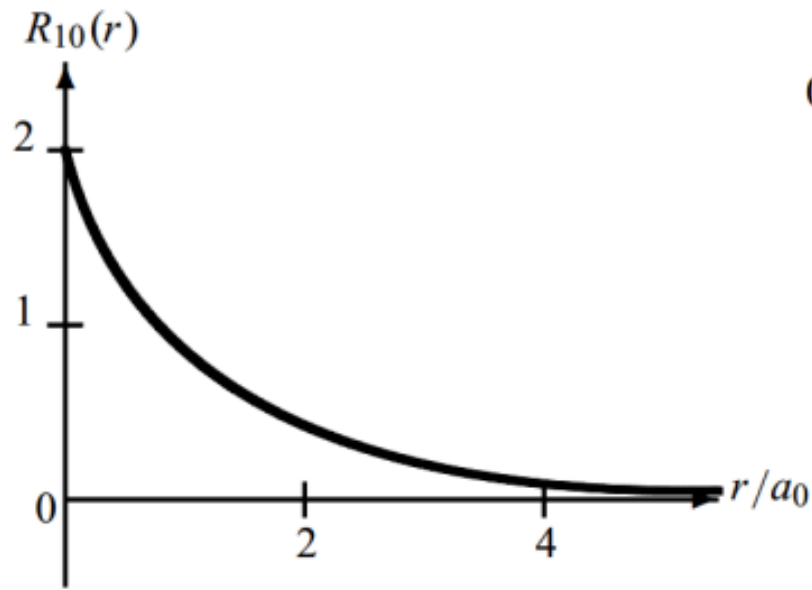
Radial eigenfunctions of hydrogenic atom

$$R_{30}(r) = 2 \left(\frac{Z}{3a_{\mu}} \right)^{3/2} \left(1 - \frac{2Zr}{3a_{\mu}} + \frac{2Z^2r^2}{27a_{\mu}^2} \right) e^{-Zr/3a_{\mu}}$$

$$R_{31}(r) = \frac{4\sqrt{2}}{9} \left(\frac{Z}{3a_{\mu}} \right)^{3/2} \left(1 - \frac{Zr}{6a_{\mu}} \right) \left(\frac{Zr}{a_{\mu}} \right) e^{-Zr/3a_{\mu}}$$

$$R_{32}(r) = \frac{4}{27\sqrt{10}} \left(\frac{Z}{3a_{\mu}} \right)^{3/2} \left(\frac{Zr}{a_{\mu}} \right)^2 e^{-Zr/3a_{\mu}}$$





The hydrogenic wave function

The solutions of the hydrogenic Schrödinger equation in spherical polar coordinates can now be written in full

$$\psi_{nlm}(r, \theta, \phi) = R_{nl}(r)Y_{lm}(\theta, \phi)$$

where $n = 1, 2, 3, \dots$ is the principle quantum number, $l = 0, 1, 2, \dots, n-1$ is the orbital angular momentum quantum number and $m = 0, \pm 1, \pm 2 \dots \pm l$ is the magnetic quantum number.



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